One-point PDF closure model applied to attached and separated flows

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Overview

- 1. Background on PDF modeling
- 2. Solution method
- 3. Preliminary results for a plane jet
- 4. Conclusions and Outlook





RANS modeling

- Start with exact NS-equations
- Introduce Reynolds averaging operator $\langle \cdot \rangle$
- Derive moment equations $\langle U_i \rangle$, $\langle u_i u_j \rangle$, ...
- Introduce models to "close" at a given moment level $\frac{D\langle u_i u_j \rangle}{Dt} = F(\langle U_i \rangle, \langle u_i u_j \rangle, \varepsilon)$

PDF modeling

- Define one-point velocity probability density function $f_e(V; x, t)$
- Start from exact NS-equations and derive exact (but unclosed) evolution equation for the PDF $\frac{Df_e}{Dt} = RHS$
- Introduce a closure model for the PDF equation and a time-scale model
- All moments are closed i.e. $\langle u_i u_j u_k \rangle$ is also closed
- Solving the closed PDF equation (7-dimensional) is not trivial!



Motivation for PDF modeling approach

- In turbulent reacting flows the chemical source terms and turbulent fluxes are treated exactly if $f(V, \psi; x, t)$ is considered
- The triple correlations $\langle u_i u_j u_k \rangle$ are closed -> possibly improved results in flows where turbulent transport is important (i.e. swirling jets)
- If we can optimize the closure model, we have directly obtained the best possible RANS model!
- New physical insights could be gained by considering Lagrangian statistics
- Could be taken to the two-point level -> Joint velocity gradient PDF

Main disadvantage: Transported PDF solution is typically 5-20 times more expensive!

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Incompressible NS – equations (constant properties) $\frac{\partial U_i}{\partial x_i} = 0$ $\frac{DU_i}{Dt} = \frac{\partial U_i}{\partial t} + U_j \frac{\partial U_i}{\partial x_j} = A_i = -\frac{1}{\rho} \frac{\partial p}{\partial x_i} + \nu \Delta^2 U_i$

Exact PDF evolution equation (derivation see e.g. Pope, "Turbulent Flows")

$$\begin{aligned} \frac{\partial f_e}{\partial t} + V_i \frac{\partial f_e}{\partial x_i} &= -\frac{\partial}{\partial V_i} \left[f_e \langle A_i | \mathbf{V} \rangle \right] = -\frac{\partial}{\partial V_i} \left[f_e \left(-\frac{1}{\rho} \frac{\partial \langle p \rangle}{\partial x_i} + \nu \Delta^2 \langle U_i \rangle \right) \right] \\ &- \frac{\partial}{\partial V_i} \left[f_e \left\langle -\frac{1}{\rho} \frac{\partial p'}{\partial x_i} \right| \mathbf{V} \right\rangle + f_e \langle \nu \Delta^2 u_i | \mathbf{V} \rangle \right] \quad \underbrace{M_i} \end{aligned}$$

Modeled PDF evolution equation

$$\frac{\partial f}{\partial t} + V_i \frac{\partial f}{\partial x_i} = -\frac{\partial}{\partial V_i} [fM_i] - \frac{\partial}{\partial V_i} [fG_{ij} (V_j - \langle U_j \rangle)] + \frac{1}{2} C_0 \varepsilon \frac{\partial^2 f}{\partial V_i \partial V_i}$$

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Modeled PDF evolution equation

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Can we "measure" the tensor G_{ij} ?

1) Pope, JFM 2014

- With moderate assumptions one finds $G_{ij} \sim \langle a_i u_k \rangle \langle u_k u_i \rangle^{-1}$
- The fluctuating velocity-acceleration correlation (a_iu_k) arises from large scale turbulent motions -> could be measured or taken from well resolved LES

2) Work done by Stefan Heinz & Grigory Sarnitsky at UWYO

• Find G_{ij} from Lagrangian DNS data using

$$\lim_{\tau\to 0}\frac{1}{\tau}\Big\langle \left[U_i'(t+\tau)-U_i'(t)\right]U_j'(t)\Big| \boldsymbol{X}(t)=\boldsymbol{x}\Big\rangle = \widetilde{G}_{ik}(\boldsymbol{x},t)\langle u_k'u_j'\rangle(\boldsymbol{x},t)$$

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Basic PDF closure model: Simplified Langevin model (SLM)

$$\frac{\partial f}{\partial t} + V_i \frac{\partial f}{\partial x_i} = -\frac{\partial}{\partial V_i} [fM_i] - \frac{\partial}{\partial V_i} [fG_{ij} (V_j - \langle U_j \rangle)] + \frac{1}{2} C_0 \varepsilon \frac{\partial^2 f}{\partial V_i \partial V_i}$$
$$G_{ij} = -\left(\frac{1}{2} + \frac{3}{4} C_0\right) \frac{\varepsilon}{k} \delta_{ij}$$

Corresponds to Rotta's return to isotropy model in RSM!

Numerical Solution

- PDF equation is 7-dimensional
- Important constraints need to be satisfied numerically f(V; x, t) > 0, and $\int f_e(V; x, t) dV = 1$

Best solved with a Monte Carlo method!

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"Particle" Monte Carlo method

SLM is a stochastic model for fluid particles which are at $X^*(t)$





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Hybrid FV – Monte Carlo solution method

Statistical error of Monte Carlo Method decreases as $\frac{1}{\sqrt{N_n}}$

-> To reduce the effects of the statistical and bias errors we use a hybrid FV – Monte Carlo method to solve for the joint velocity – turbulence frequency PDF

- Consider a FV mesh with Nc = 50-100 particles per cell
- Solve mean velocity $\langle U_i \rangle$ and pressure equations with FV method and Reynolds stress taken from particles (particle in cell average + time averaging)
- Solve particle equation with a simple Euler-Maruyama method

$$u_{i}^{*}(t + \Delta t) = u_{i}^{*}(t) + \frac{\partial \langle u_{i}u_{j} \rangle}{\partial x_{j}} \Delta t - u_{j}^{*}(t) \frac{\partial \langle U_{i} \rangle}{\partial x_{j}} - \left(\frac{1}{2} + \frac{3}{4}C_{0}\right) \Omega u_{j}^{*}(t) + (C_{0} k \Omega \Delta t)^{1/2} \xi_{i}$$

$$\omega^{*}(t + \Delta t) = \omega^{*}(t) - C_{\omega,3}(\omega^{*} - \langle \omega \rangle) \Omega \Delta t - S_{\omega} \Omega \omega^{*} \Delta t + \left(2C_{\omega,3}C_{\omega,4} \langle \omega \rangle \omega^{*} \Omega \Delta t\right)^{1/2} \xi'$$

$$X_{i}^{*}(t + \Delta t) = X_{i}^{*}(t) + \left(\langle U_{i} \rangle + u_{i}^{*}(t + \Delta t)\right) dt$$

$$\Omega = C_{\Omega} \langle \omega^{*} | \omega^{*} > \langle \omega \rangle \rangle$$



 $\langle uu \rangle$



Preliminary results: Plane jet at Re = 22,000

"Simultaneous measurement of all three velocity components and pressure in a plane jet" Terashima et al, Meas. Sci. Technol. 25 (2014)





- Measured velocity and joint pressure-velocity statistics (one-point)
- Performed LES to "Validate" the measurements

Simulations

- Hybrid FV- Monte Carlo solver implemented in OpenFOAM (based on pisoFOAM)
- Implementation suitable for unstructured grids and fully parallelized
- 50d x 20d domain with 240x110 grid points
- Np = 100 particles per cell, time averaging over
 200 time steps
- For comparison we use a Rotta's RSM model
 with Daily-Harlow turbulent diffusion and standard ε-equation coefficients (RottaDH)
 PDF solution is based on a "frozen" mean velocity field from RottaDH solution

Preliminary results: Plane jet at Re = 22,000

PDF: $\langle uu \rangle$

RottaDH RSM: $\langle uu \rangle$







Preliminary results: Plane jet at Re = 22,000 x/h = 20

PDF RSM

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PDF RSM

3

2

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2

3



Preliminary results: Plane jet at Re = 22,000 x/h = 20







0.4

0.2

0

-3

-2

-1

0

y/b

1

2

3

Conclusions and Outlook

- PDF closure modeling provides a "complete" momentclosure
- Importance of higher order moments such as the turbulent transport could be evaluated and new models for RSM approaches proposed
- A basic hybrid FV Monet Carlo solution method has been implemented in OpenFOAM (unstructured grids, fully parallelized with encouraging strong scaling)
- More debugging required
- Analyze k-budget for the plane jet also for different G_{ij} models

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